

Third order Maximum-Principle-Satisfying Direct DG methods for convection diffusion equations on unstructured triangular mesh

Zheng Chen

chenz1@ornl.gov

CAM Semiar
University of Tennessee
March 2, 2016

Object: design high order Maximum-Principle-Satisfying DG method

$$u_t + \nabla \cdot F(u) - \nabla \cdot (A(u)\nabla u) = 0$$

- ▶ PDE (continuous): The solution satisfies the maximum principle:

$$m \leq u(x, y, t) \leq M, \quad (x, y) \in \Omega, t > 0,$$

where, m and M are the minimum and maximum of the initial and boundary data.

- ▶ Schemes (discrete): We expect the numerical solution to satisfy the discrete maximum principle at all time level t^n :

$$m \leq u_K^n(x, y) \leq M, \quad (x, y) \in K \in \mathcal{T}_h \subseteq \Omega.$$

- ▶ **Strong stability** result in the L^∞ sense
- ▶ Consistency to physical meanings (e.g. preserve **positivity** of density profile)

M-P-S high order DG methods for hyperbolic PDEs

$$u_t + \nabla \cdot F(u) = 0 \quad \implies \quad m \leq u_K^n(x, y) \leq M$$

For convection term:

M-P-S DGM on conservation laws

- ▶ X. Zhang and C.-W. Shu (2010a). "On maximum-principle-satisfying high order schemes for scalar conservation laws". In: *J. Comput. Phys.*
- ▶ X. Zhang and C.-W. Shu (2011). "Maximum-principle-satisfying and positivity-preserving high-order schemes for conservation laws: survey and new developments". In: *Proc. R. Soc. A*
- ▶ X. Zhang, Y. Xia, and C.-W. Shu (2012). "Maximum-principle-satisfying and positivity-preserving high order discontinuous Galerkin schemes for conservation laws on triangular meshes". In: *J. Sci. Comput.*

$$u_t + \nabla \cdot F(u) - \nabla \cdot (A(u)\nabla u) = 0 \quad \implies \quad m \leq u_K^n(x, y) \leq M$$

For diffusion term:

- ▶ X. Zhang, Y. Zhang and C.-W. Shu (2013). "Maximum-principle-satisfying second order discontinuous Galerkin schemes for convection-diffusion equations on triangular meshes". In: *J. Comput. Phys.*
- ▶ Z. Chen, H. Huang, and J. Yan (2016). "Third order maximum-principle-satisfying direct discontinuous Galerkin methods for time dependent convection diffusion equations on unstructured triangular meshes". In: *Journal of Computational Physics* 308, pp. 198–217

Outline

1 Reviews

- Discontinuous Galerkin methods
- General framework of M-P-S schemes
- Introduction to Direct DG method and its variations

2 M-P-S DDG schemes for triangular mesh

- Methods on rectangular mesh
- New ideas designed for unstructured mesh
- Extension to nonlinear diffusion equations
- Algorithm

3 Numerical examples

4 Conclusions and future works

Joint work with: Jue Yan (Iowa State University)
Hongying Huang (Zhejiang Ocean University, China)

Outline

1 Reviews

- Discontinuous Galerkin methods
- General framework of M-P-S schemes
- Introduction to Direct DG method and its variations

2 M-P-S DDG schemes for triangular mesh

- Methods on rectangular mesh
- New ideas designed for unstructured mesh
- Extension to nonlinear diffusion equations
- Algorithm

3 Numerical examples

4 Conclusions and future works

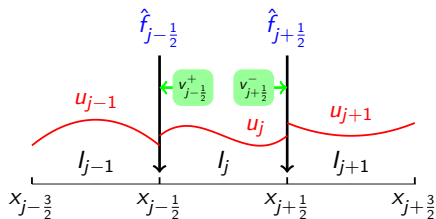
Review of discontinuous Galerkin (DG) methods

$$\boxed{u_t + f(u)_x = 0} \implies \boxed{\int_{I_j} u_t v \, dx + \int_{I_j} f(u)_x v \, dx = 0}$$

Integration by part \implies

$$\int_{I_j} u_t v \, dx - \int_{I_j} f(u) v_x \, dx + [f(u)v] \Big|_{x_{j-\frac{1}{2}}}^{x_{j+\frac{1}{2}}} = 0$$

DG solution:



$$\int_{I_j} u_t v \, dx - \int_{I_j} f(u) v_x \, dx + \hat{f}_{j+\frac{1}{2}} v_{j+\frac{1}{2}}^- - \hat{f}_{j-\frac{1}{2}} v_{j-\frac{1}{2}}^+ = 0$$

Outline

1 Reviews

- Discontinuous Galerkin methods
- **General framework of M-P-S schemes**
- Introduction to Direct DG method and its variations

2 M-P-S DDG schemes for triangular mesh

- Methods on rectangular mesh
- New ideas designed for unstructured mesh
- Extension to nonlinear diffusion equations
- Algorithm

3 Numerical examples

4 Conclusions and future works

Major steps to prove M-P-S schemes

Assumption:

$$|u(x, t^n) - u_j^n(x)| \leq C\Delta x^3 \text{ and } m \leq u_j^n(x) \leq M, \forall x \in I_j$$

- ▶ **Step 1:** evolve in time once and estimate the solution **average**

$$\bar{u}_j^{n+1} = \frac{1}{\Delta x_j} \int_{I_j} u_j^{n+1}(x) dx$$

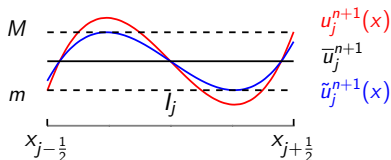
$$m \leq \bar{u}_j^{n+1} \leq M$$

← difficult!

- ▶ **Step 2:** modify the solution by adding a limiter to make sure

$$m \leq \tilde{u}_j^{n+1}(x) \leq M$$

← easy



$$\tilde{u}_j^{n+1}(x) = \theta(u_j^{n+1}(x) - \bar{u}_j^{n+1}) + \bar{u}_j^{n+1}$$
$$\theta = \min \left\{ 1, \left| \frac{M - \bar{u}_j^{n+1}}{M_j - \bar{u}_j^{n+1}} \right|, \left| \frac{m - \bar{u}_j^{n+1}}{m_j - \bar{u}_j^{n+1}} \right| \right\}$$

Difficulty to obtain high order M-P-S for diffusion equation

Hyperbolic

$$u_t + f(u)_x = 0$$



$$\int_{I_j} u_t v \, dx - \int_{I_j} f(u) v_x \, dx =$$
$$-\hat{f}_{j+\frac{1}{2}} v_{j+\frac{1}{2}}^- + \hat{f}_{j-\frac{1}{2}} v_{j-\frac{1}{2}}^+$$



$$\bar{u}_j^{n+1} = \bar{u}_j^n + \frac{\Delta t}{\Delta x} \left(\hat{f}_{j+\frac{1}{2}} - \hat{f}_{j-\frac{1}{2}} \right)$$

solution values
on cell boundaries

Parabolic

$$u_t = u_{xx}$$



$$\int_{I_j} u_t v \, dx + \int_{I_j} u_x v_x \, dx =$$
$$(\widehat{u_x})_{j+\frac{1}{2}} v_{j+\frac{1}{2}}^- - (\widehat{u_x})_{j-\frac{1}{2}} v_{j-\frac{1}{2}}^+$$



$$\bar{u}_j^{n+1} = \bar{u}_j^n + \frac{\Delta t}{\Delta x} \left((\widehat{u_x})_{j+\frac{1}{2}} - (\widehat{u_x})_{j-\frac{1}{2}} \right)$$

solution derivative values
on cell boundaries

Outline

1 Reviews

- Discontinuous Galerkin methods
- General framework of M-P-S schemes
- Introduction to Direct DG method and its variations

2 M-P-S DDG schemes for triangular mesh

- Methods on rectangular mesh
- New ideas designed for unstructured mesh
- Extension to nonlinear diffusion equations
- Algorithm

3 Numerical examples

4 Conclusions and future works

Direct DG (DDG) method as a diffusion solver

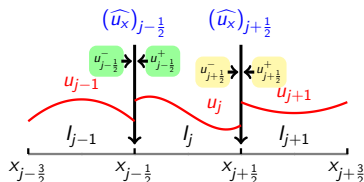
Parabolic

$$u_t = u_{xx} \implies \int_{I_j} u_t v \, dx + \int_{I_j} u_x v_x \, dx - (\widehat{u_x})_{j+\frac{1}{2}} v_{j+\frac{1}{2}}^- + (\widehat{u_x})_{j-\frac{1}{2}} v_{j-\frac{1}{2}}^+ = 0$$

At cell boundaries:

$$[u] = u^+ - u^-$$

$$\bar{u} = \frac{u^+ + u^-}{2}$$



$$\widehat{u_x} \Big|_{x_{j\pm\frac{1}{2}}} = \beta_0 \frac{[u]}{\Delta x} + \bar{u_x} + \beta_1 \Delta x [u_{xx}] + \beta_2 \Delta x^3 [u_{xxxx}] + \dots^*$$

*H. Liu and J. Yan (2009). "The Direct Discontinuous Galerkin (DDG) Methods for Diffusion Problems". In: *SIAM J. Numer. Anal.*

DDG method with interface correction

$$\int_{I_j} u_t v \, dx + \int_{I_j} u_x v_x \, dx - (\widehat{u}_x v) \Big|_{x_{j-\frac{1}{2}}}^{x_{j+\frac{1}{2}}} + \boxed{\sigma \left([u](\widetilde{v}_x)_{j+\frac{1}{2}} + [u](\widetilde{v}_x)_{j-\frac{1}{2}} \right)} = 0$$

- ▶ DDG: $\sigma = 0$, $\widehat{u}_x = \beta_0 \frac{[u]}{\Delta x} + \overline{u}_x + \beta_1 \Delta x [u_{xx}] + \beta_3 \Delta x^3 [u_{xxxx}] + \dots$
- ▶ DDG with interface correction: $\sigma = 1$,

$$\begin{cases} \widehat{u}_x = \beta_0 \frac{[u]}{\Delta x} + \overline{u}_x + \beta_1 \Delta x [u_{xx}] \\ \widetilde{v}_x = \overline{v}_x \end{cases}$$

DDG method with interface correction

$$\int_{I_j} u_t v \, dx + \int_{I_j} u_x \cancel{v_x} \, dx - (\widehat{u}_x v) \Big|_{x_{j-\frac{1}{2}}}^{x_{j+\frac{1}{2}}} + \sigma \left([u] \cancel{(\widetilde{v}_x)}_{j+\frac{1}{2}} + [u] \cancel{(\widetilde{v}_x)}_{j-\frac{1}{2}} \right) = 0$$

- ▶ DDG: $\sigma = 0$, $\widehat{u}_x = \beta_0 \frac{[u]}{\Delta x} + \overline{u_x} + \beta_1 \Delta x [u_{xx}] + \beta_3 \Delta x^3 [u_{xxxx}] + \dots$
- ▶ DDG with interface correction: $\sigma = 1$,

$$\begin{cases} \widehat{u}_x = \beta_0 \frac{[u]}{\Delta x} + \overline{u_x} + \beta_1 \Delta x [u_{xx}] \\ \widetilde{v}_x = \overline{v_x} \end{cases}$$

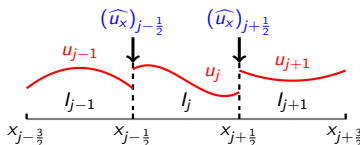
Solution average evolution

Take $v = 1$

$$\frac{d}{dt} \int_{I_j} u \, dx = (\widehat{u_x}) \Big|_{x_{j-\frac{1}{2}}}^{x_{j+\frac{1}{2}}} \implies \frac{d}{dt} \bar{u} = \frac{1}{\Delta x} \left\{ (\widehat{u_x})_{j+\frac{1}{2}} - (\widehat{u_x})_{j-\frac{1}{2}} \right\}$$

Euler forward in time evolution

$$\begin{aligned} \bar{u}_j^{n+1} &= \bar{u}_j^n + \frac{\Delta t}{\Delta x} \left\{ (\widehat{u_x})_{j+\frac{1}{2}} - (\widehat{u_x})_{j-\frac{1}{2}} \right\} \\ &= H(u_{j-1}^n(x), u_j^n(x), u_{j+1}^n(x)) \end{aligned}$$



Goal

$$m \leq u_{j-1}^n(x), u_j^n(x), u_{j+1}^n(x) \leq M \implies m \leq \bar{u}_j^{n+1} \leq M$$

Outline

1 Reviews

- Discontinuous Galerkin methods
- General framework of M-P-S schemes
- Introduction to Direct DG method and its variations

2 M-P-S DDG schemes for triangular mesh

- **Methods on rectangular mesh**
- New ideas designed for unstructured mesh
- Extension to nonlinear diffusion equations
- Algorithm

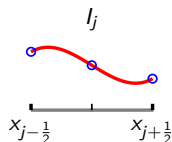
3 Numerical examples

4 Conclusions and future works

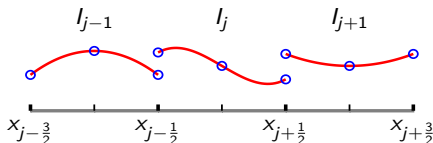
Previous method to bound solution average

Previous idea

Pick $u_{j-\frac{1}{2}}^+$, u_j , $u_{j+\frac{1}{2}}^-$ as degrees of freedom and rewrite the quadratic polynomial $u_j^u(x)$ in the Lagrange format.

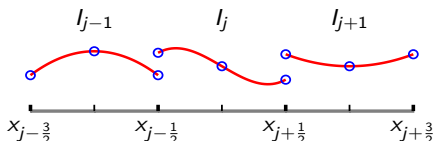


$$\begin{aligned}\bar{u}_j^{n+1} &= H(u_{j-1}^n(x), u_j^n(x), u_{j+1}^n(x)) \\ &= H(u_{j-\frac{3}{2}}^+, u_{j-1}, u_{j-\frac{1}{2}}^-, u_{j-\frac{1}{2}}^+, u_j, u_{j+\frac{1}{2}}^-, u_{j+\frac{1}{2}}^+, u_{j+1}, u_{j+\frac{3}{2}}^-)\end{aligned}$$



M-P-S 3rd order DDG schemes

$$\begin{aligned}\bar{u}_j^{n+1} &= H(u_{j-\frac{3}{2}}^+, u_{j-1}, u_{j-\frac{1}{2}}^-, u_{j-\frac{1}{2}}^+, u_j, u_{j+\frac{1}{2}}^-, u_{j+\frac{1}{2}}^+, u_{j+1}, u_{j+\frac{3}{2}}^-) \\ &= H(\uparrow, \uparrow, \uparrow, \uparrow, \uparrow, \uparrow, \uparrow, \uparrow, \uparrow) \text{ convex combination}\end{aligned}$$



Given $m \leq u_j^n(x) \leq M$, we have

$$m = H(m, \dots, m) \leq \bar{u}_j^{n+1} \leq H(M, \dots, M) = M$$

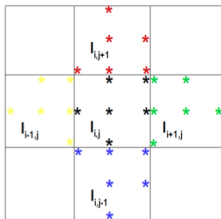
Conditions needed:

$$\beta_0 \geq \frac{3}{2} \quad \frac{1}{8} \leq \beta_1 \leq \frac{1}{4}$$

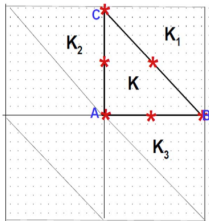
Same results holds for SSP
Runge Kutta methods.

$$\widehat{u}_x = \beta_0 \frac{[u]}{\Delta x} + \bar{u}_x + \beta_1 \Delta x [u_{xx}]$$

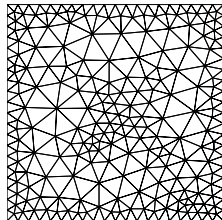
Difficulty for unstructured triangular mesh



(a) rectangular mesh



(b) triangular mesh



(c) unstructured triangular mesh

For any shape triangle, it is very hard to identify the **6 degrees of freedom** to rewrite the quadratic polynomial $u_K^n(x, y)$ to bound

$$m \leq \bar{u}_K^{n+1} \leq M$$

Outline

1 Reviews

- Discontinuous Galerkin methods
- General framework of M-P-S schemes
- Introduction to Direct DG method and its variations

2 M-P-S DDG schemes for triangular mesh

- Methods on rectangular mesh
- New ideas designed for unstructured mesh
- Extension to nonlinear diffusion equations
- Algorithm

3 Numerical examples

4 Conclusions and future works

New idea for triangular mesh

$$u_t - \Delta u = 0$$

DDG schemes:

$$\int_K u_t v \, dx dy + \int_K \nabla u \cdot \nabla v \, dx dy - \int_{\partial K} \widehat{u}_n v \, ds = 0$$

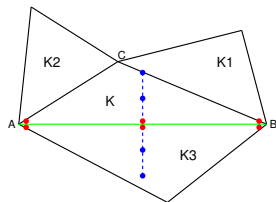
Numerical flux:

$$\widehat{u}_n = \beta_0 \frac{[u]}{h_K} + \bar{u}_n + \beta_1 h_K [u_{nn}]$$

Solution average evolution:

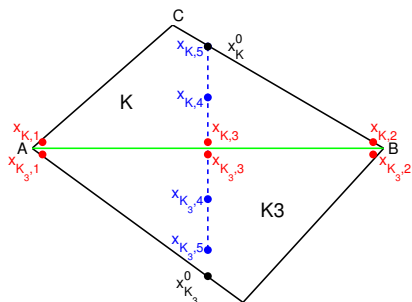
$$\bar{u}_K^{n+1} = \bar{u}_K^n + \frac{\Delta t}{\text{area}(K)} \int_{\partial K} \widehat{u}_n \, ds$$

New idea: calculate $\int_{\partial K} \widehat{u}_n v \, ds$ directly!



New idea for triangular mesh

$$\bar{u}_K^{n+1} = \bar{u}_K^n + \frac{\Delta t}{\text{area}(K)} \left\{ \int_{AB} \hat{u}_n ds + \int_{BC} \hat{u}_n ds + \int_{CA} \hat{u}_n ds \right\}$$

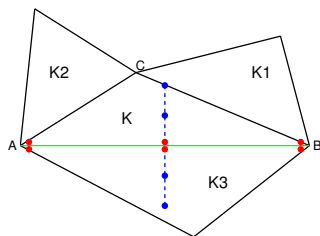


Let's use edge AB shared by elements K and K_3 to illustrate the solution points selected to calculate

$$\int_{AB} \hat{u}_n ds = \frac{\beta_0}{h_K} \int_{AB} [u] ds + \int_{AB} \bar{u}_n ds + \beta_1 h_K \int_{AB} [u_{nn}] ds.$$

$u_K^n(x, y)$ is a **quadratic** polynomial.

3rd order M-P-S DDG schemes



$$\bar{u}_K^{n+1} = H(u_K^n(x, y), u_{K_1}^n(x, y), u_{K_2}^n(x, y), u_{K_3}^n(x, y)) = H(\uparrow, \uparrow, \uparrow, \uparrow)$$

Theorem

Given $m \leq u_K^n(x, y) \leq M$ for all K , we have $m \leq \bar{u}_K^{n+1} \leq M$ provided,

$$\beta_0 \geq \frac{9}{4} - 6\beta_1, \quad \frac{1}{8} \leq \beta_1 \leq \frac{1}{4}, \quad \lambda = \frac{\Delta t}{\text{area}(K)} \leq C(\beta_0, \beta_1, T_h).$$

Outline

1 Reviews

- Discontinuous Galerkin methods
- General framework of M-P-S schemes
- Introduction to Direct DG method and its variations

2 M-P-S DDG schemes for triangular mesh

- Methods on rectangular mesh
- New ideas designed for unstructured mesh
- **Extension to nonlinear diffusion equations**
- Algorithm

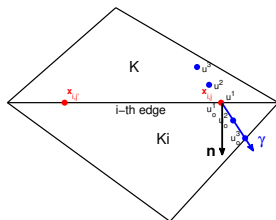
3 Numerical examples

4 Conclusions and future works

Nonlinear diffusion on triangular mesh

$$u_t - \nabla \cdot (A(u)\nabla u) = 0$$

$$\bar{u}_K^{n+1} = \bar{u}_K^n + \frac{\Delta t}{\text{area}(K)} \int_{\partial K} (A(\widehat{u})\widehat{\nabla u} \cdot \mathbf{n}) \, ds$$



We have

$$A(u)\nabla u \cdot \mathbf{n} = \nabla u \cdot \gamma = \|\gamma\| u_\gamma \quad \text{with} \quad \gamma = A^T(u)\mathbf{n}$$

Diffusion matrix $A(u)$ is positive definite $\implies \gamma \cdot \mathbf{n} = \mathbf{n}^T A(u)\mathbf{n} > 0$

Numerical flux:

$$A(\widehat{u})\widehat{\nabla u} \cdot \mathbf{n} = \|\gamma\| \widehat{u}_\gamma = \|\gamma\| \left\{ \beta_0 \frac{[u]}{h_K} + \bar{u}_\gamma + \beta_1 h_K [u_\gamma \gamma] \right\}$$

Outline

1 Reviews

- Discontinuous Galerkin methods
- General framework of M-P-S schemes
- Introduction to Direct DG method and its variations

2 M-P-S DDG schemes for triangular mesh

- Methods on rectangular mesh
- New ideas designed for unstructured mesh
- Extension to nonlinear diffusion equations
- **Algorithm**

3 Numerical examples

4 Conclusions and future works

2-D M-P-S limiter

$$\tilde{u}_K^n(x, y) = \theta(u_K^n(x, y) - \bar{u}_K^n) + \bar{u}_K^n, \quad \theta = \min \left\{ 1, \left| \frac{M - \bar{u}_K^n}{M_K - \bar{u}_K^n} \right|, \left| \frac{m - \bar{u}_K^n}{m_K - \bar{u}_K^n} \right| \right\},$$

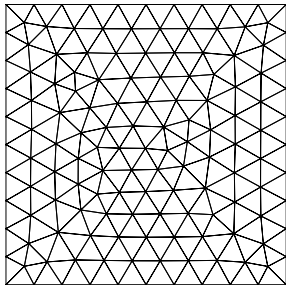
with

$$M_K = \max_{(x,y) \in K} u_K^n(x, y), \quad m_K = \min_{(x,y) \in K} u_K^n(x, y).$$

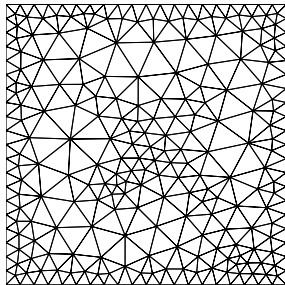
Algorithm

1. At time level t^n , apply M-P-S limiter to $u_K^n(x, y)$ and obtain $\tilde{u}_K^n(x, y)$
2. Apply DDG scheme to $\tilde{u}_K^n(x, y)$ and evolve in time with SSP RK method and march forward the solution to the next time level $u_K^{n+1}(x, y)$

Mesh



(a) Triangular mesh



(b) Mesh with obtuse triangles

Accuracy Tests

$$u_t - \epsilon \Delta u = 0 \quad \text{with} \quad \epsilon = 1$$

$$\text{exact solution: } u(x, y, t) = e^{-8\pi^2 \epsilon t} \sin(2\pi(x + y))$$

order h	no M-P-S limiter				with M-P-S limiter			
	L^2 error	L^∞ error	$u_{\min} - ue_{\min}$	$u_{\max} - ue_{\max}$	L^2 error	L^∞ error	$u_{\min} - ue_{\min}$	$u_{\max} - ue_{\max}$
0.117			-3.91e-03	1.28e-03			0	0
0.0587	2.64	3.10	-2.40e-04	1.95e-04	2.71	3.10	0	0
0.0293	2.90	3.01	-1.49e-05	2.60e-05	2.91	3.01	0	0
0.0147	2.99	2.98	-9.61e-07	3.27e-06	2.99	2.98	0	0
0.00733	2.99	2.99	-5.82e-08	4.11e-07	2.99	2.99	0	0

Table: Accuracy table on triangular mesh 1(a) at $t = 0.0001$.

order h	no M-P-S limiter				with M-P-S limiter			
	L^2 error	L^∞ error	$u_{\min} - ue_{\min}$	$u_{\max} - ue_{\max}$	L^2 error	L^∞ error	$u_{\min} - ue_{\min}$	$u_{\max} - ue_{\max}$
0.148			-1.17e-02	9.70e-03			0	0
0.0741	2.74	3.00	-7.86e-04	5.49e-05	3.11	2.85	0	0
0.0371	2.87	3.06	-2.66e-05	4.24e-05	2.90	3.22	0	0
0.0185	2.97	3.02	-1.12e-06	9.22e-07	2.97	3.02	0	0
0.00927	2.99	3.01	-7.17e-08	1.33e-08	2.99	3.01	0	0

Table: Accuracy table on unstructured mesh 1(b) with obtuse triangles at 0.0001.

2-D porous medium equations

$$u_t - \Delta(u^2) = 0$$

Positivity-preserving $m = 0$

2-D porous medium equations

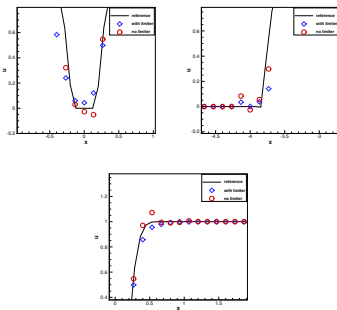
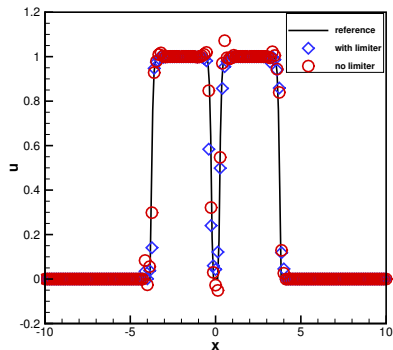


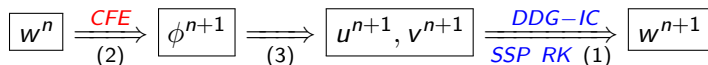
Figure: The cut of the DDG interface correction solution along line $x + y = 0$ at $t = 0.005$. Red circle symbol: no M-P-S limiter. Blue diamond symbol: M-P-S limiter applied.

Incompressible Navier-Stokes equations

vorticity stream-function formulation

$$\left\{ \begin{array}{l} w_t + (uw)_x + (vw)_y = \frac{1}{Re} \nabla^2 w \\ \Delta \phi = w \\ \langle u, v \rangle = \langle -\phi_y, \phi_x \rangle \end{array} \right. \quad \begin{array}{l} (1) \\ (2) \\ (3) \end{array}$$

We adopt Liu and Shu's method[†] that couple stream function ϕ with P^2 continuous finite element and DG for vorticity w .



We carry out two tests: accuracy check and vortex patch problem.

[†]J.-G. Liu and C.-W. Shu (2000). "A high-order discontinuous Galerkin method for 2D incompressible flows". In: *J. Comput. Phys.*

Accuracy check

$$w(x, y, t) = -2 \sin(x) \sin(y) e^{-\frac{2t}{Re}}$$

order $h/2\pi$	no M-P-S limiter				with M-P-S limiter			
	L^2 error	L^∞ error	$u_{\min} - ue_{\min}$	$u_{\max} - ue_{\max}$	L^2 error	L^∞ error	$u_{\min} - ue_{\min}$	$u_{\max} - ue_{\max}$
0.117			-8.17e-03	7.72e-03			0	0
0.0587	2.82	3.16	-3.64e-04	6.41e-04	2.88	3.18	0	0
0.0293	2.89	2.82	-6.40e-05	-4.50e-06	2.89	2.82	0	-4.52e-06
0.0147	3.00	3.51	-1.29e-06	3.62e-06	3.00	3.51	0	0
0.00733	2.96	2.66	-1.53e-07	1.95e-07	2.96	2.66	0	0

Table: Mesh 1(a), final time $T = 0.1$

order $h/2\pi$	no M-P-S limiter				with M-P-S limiter			
	L^2 error	L^∞ error	$u_{\min} - ue_{\min}$	$u_{\max} - ue_{\max}$	L^2 error	L^∞ error	$u_{\min} - ue_{\min}$	$u_{\max} - ue_{\max}$
0.148			-8.40e-04	-7.18e-05			0	-7.18e-05
0.0741	2.93	3.03	3.56e-05	-6.40e-05	2.93	3.03	4.15e-05	-6.40e-05
0.0371	2.86	2.90	3.36e-06	1.71e-06	2.86	2.90	3.36e-06	0
0.0185	2.99	3.07	1.79e-07	1.43e-07	2.99	3.07	1.79e-07	0

Table: Mesh 1(b) with obtuse triangles, final time $T = 0.1$

Vortex patch problem

$$\begin{cases} w_t + (uw)_x + (vw)_y = \frac{1}{Re} \nabla w \\ \Delta \phi = w, \quad \langle u, v \rangle = \langle -\phi_y, \phi_x \rangle \end{cases}$$

$$w_0(x, y) = \begin{cases} -1, & (x, y) \in [\frac{\pi}{2}, \frac{3\pi}{2}] \times [\frac{\pi}{4}, \frac{3\pi}{4}], \\ 1, & (x, y) \in [\frac{\pi}{2}, \frac{3\pi}{2}] \times [\frac{5\pi}{4}, \frac{7\pi}{4}], \\ 0, & \textit{otherwise.} \end{cases}$$

Vortex patch $Re = 100$

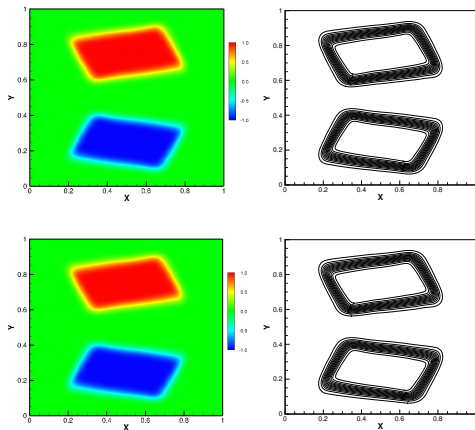


Figure: Contours of the solutions, $Re = 100$, at $t = 1$. Top: no M-P-S limiter ($w_{\min} = -1.0011$, $w_{\max} = 1.0008$); Bottom: add M-P-S limiter ($w_{\min} = -1$, $w_{\max} = 1$). 30 equally spaced contour lines are plotted.

Vortex patch $Re = 10000$

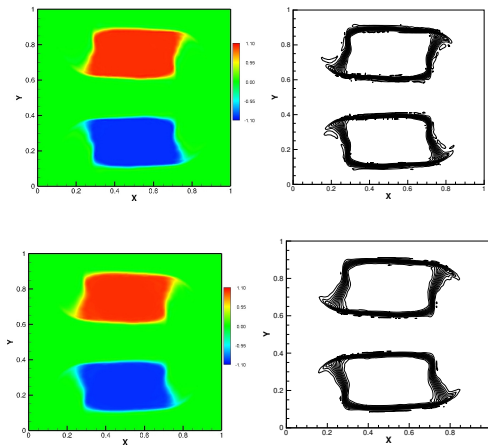


Figure: Contours of the solutions, $Re = 10000$ at $t = 5$. Top: no M-P-S limiter ($w_{\min} = -1.1211$, $w_{\max} = 1.1508$); Bottom: add M-P-S limiter ($w_{\min} = -1$, $w_{\max} = 1$). 30 equally spaced contour lines are plotted.

Vortex patch $Re = 10000$

Time evolutions of numerical vorticity w from DDG-IC schemes with M-P-S limiter, up to $t = 5$.

Vortex patch

$h/2\pi$	$Re = 100$				$Re = 10000$			
	no limiter		with limiter		no limiter		with limiter	
	$w_{\min} - we_{\min}$	$w_{\max} - we_{\max}$	$w_{\min} - we_{\min}$	$w_{\max} - we_{\max}$	$w_{\min} - we_{\min}$	$w_{\max} - we_{\max}$	$w_{\min} - we_{\min}$	$w_{\max} - we_{\max}$
0.117	-6.96e-01	6.87e-01	0	0	-8.62e-01	8.62e-01	0	0
0.0587	-2.72e-01	2.66e-01	0	0	-6.14e-01	7.82e-01	0	0
0.0293	-8.02e-02	4.22e-02	0	0	-5.14e-01	5.77e-01	0	0
0.0147	-3.17e-03	3.00e-03	0	0	-4.35e-01	4.79e-01	0	0
0.00733	-6.41e-04	7.41e-04	0	0	-3.52e-01	2.81e-01	0	0

Table: Maximum and minimum of the solutions, $Re = 100$ and $Re = 10000$ at $t = 0.1$.

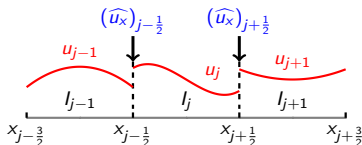
Key ingredient

$$u_t - u_{xx} = 0$$

$$\int_{I_j} u_t v \, dx + \int_{I_j} u_x v_x \, dx - (\widehat{u_x} v) \Big|_{x_{j-\frac{1}{2}}}^{x_{j+\frac{1}{2}}} + [u](\widetilde{v_x})_{j+\frac{1}{2}} + [u](\widetilde{v_x})_{j-\frac{1}{2}} = 0$$

Euler forward in time evolution ($v = 1$)

$$\begin{aligned} \bar{u}_j^{n+1} &= \bar{u}_j^n + \frac{\Delta t}{\Delta x} \left\{ (\widehat{u_x})_{j+\frac{1}{2}} - (\widehat{u_x})_{j-\frac{1}{2}} \right\} \\ &= H(u_{j-1}^n(x), u_j^n(x), u_j^{n+1}(x)) \end{aligned}$$



with $\widehat{u_x} = \beta_0 \frac{[u]}{\Delta x} + \bar{u_x} + \beta_1 \Delta x [u_{xx}]$

Degenerates to IPDG without second derivatives jump term.

Conclusions

- ▶ Sufficient conditions are given to obtain 3rd order M-P-S DDG methods for convection diffusion problems on triangular mesh
- ▶ Second derivative jump term $[u_{nn}]$ is crucial.

Future work

- ▶ Positivity-preserving for the heights of shallow water equations
- ▶ Positivity-preserving for density and pressure profiles of compressible Navier-Stokes Equations

Thank you for your attention!

Any questions?

Motivations

- ▶ **Strong stability** result in the L^∞ sense
- ▶ Consistency to physical meanings (e.g. preserve **positivity** of density profile)

Numerical challenges

- ▶ Finite difference methods: no better than 2nd order accuracy
- ▶ Classical finite element methods (2D):
 - ▶ Linear FE: mesh without obtuse elements
 - ▶ Quadratic FE: extremely restrictive assumptions on the mesh
- ▶ Little is known about high order finite volume or spectral methods, etc.

DDG method and its variations

$$\int_{I_j} u_t v \, dx + \int_{I_j} u_x v_x \, dx - (\widehat{u}_x v) \Big|_{x_{j-\frac{1}{2}}}^{x_{j+\frac{1}{2}}} + \sigma \left([u](\widetilde{v}_x)_{j+\frac{1}{2}} + [u](\widetilde{v}_x)_{j-\frac{1}{2}} \right) = 0$$

- ▶ DDG: $\sigma = 0$, $\widehat{u}_x \Big|_{x_{j\pm\frac{1}{2}}} = \beta_0 \frac{[u]}{\Delta x} + \overline{u_x} + \beta_1 \Delta x [u_{xx}] + \beta_3 \Delta x^3 [u_{xxxx}] + \dots$

- ▶ DDG with interface correction: $\sigma = 1$,

$$\begin{cases} \widehat{u}_x = \beta_0 \frac{[u]}{\Delta x} + \overline{u_x} + \beta_1 \Delta x [u_{xx}] \\ \widetilde{v}_x = \overline{v_x} \end{cases}$$

- ▶ DDG with symmetric structure: $\sigma = 1$, ($\beta_1 = 0 \rightarrow$ IPDG)

$$\begin{cases} \widehat{u}_x = \beta_{u,0} \frac{[u]}{\Delta x} + \overline{u_x} + \beta_1 \Delta x [u_{xx}] \\ \widetilde{v}_x = \beta_{v,0} \frac{[v]}{\Delta x} + \overline{v_x} + \beta_1 \Delta x [v_{xx}] \end{cases}$$

- ▶ non-symmetric DDG method: $\sigma = -1$, ($\beta_{u,0} > \beta_{v,0} \rightarrow$ NIPG)

$$\begin{cases} \widehat{u}_x = \beta_{u,0} \frac{[u]}{\Delta x} + \overline{u_x} + \beta_1 \Delta x [u_{xx}] \\ \widetilde{v}_x = \beta_{v,0} \frac{[v]}{\Delta x} + \overline{v_x} + \beta_1 \Delta x [v_{xx}] \end{cases}$$

DDG method and its variations

$$\int_{I_j} u_t v \, dx + \int_{I_j} u_x \cancel{v_x} \, dx - (\widehat{u_x} v) \Big|_{x_{j-\frac{1}{2}}}^{x_{j+\frac{1}{2}}} + \boxed{\sigma \left([u] \cancel{(\widetilde{v_x})}_{j+\frac{1}{2}} + [u] \cancel{(\widetilde{v_x})}_{j-\frac{1}{2}} \right)} = 0$$

- ▶ DDG: $\sigma = 0$, $\widehat{u_x} \Big|_{x_{j\pm\frac{1}{2}}} = \beta_0 \frac{[u]}{\Delta x} + \overline{u_x} + \beta_1 \Delta x [u_{xx}] + \beta_3 \Delta x^3 [u_{xxxx}] + \dots$

- ▶ **DDG with interface correction:** $\sigma = 1$,

$$\begin{cases} \widehat{u_x} = \beta_0 \frac{[u]}{\Delta x} + \overline{u_x} + \beta_1 \Delta x [u_{xx}] \\ \widetilde{v_x} = \overline{v_x} \end{cases}$$

- ▶ DDG with symmetric structure: $\sigma = 1$, ($\beta_1 = 0 \rightarrow$ IPDG)

$$\begin{cases} \widehat{u_x} = \beta_{u,0} \frac{[u]}{\Delta x} + \overline{u_x} + \beta_1 \Delta x [u_{xx}] \\ \widetilde{v_x} = \beta_{v,0} \frac{[v]}{\Delta x} + \overline{v_x} + \beta_1 \Delta x [v_{xx}] \end{cases}$$

- ▶ non-symmetric DDG method: $\sigma = -1$, ($\beta_{u,0} > \beta_{v,0} \rightarrow$ NIPG)

$$\begin{cases} \widehat{u_x} = \beta_{u,0} \frac{[u]}{\Delta x} + \overline{u_x} + \beta_1 \Delta x [u_{xx}] \\ \widetilde{v_x} = \beta_{v,0} \frac{[v]}{\Delta x} + \overline{v_x} + \beta_1 \Delta x [v_{xx}] \end{cases}$$

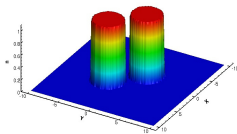
2-D porous medium equations

$$u_t - \Delta(u^2) = 0$$

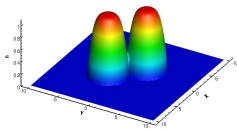
Positivity-preserving $m = 0$

2-D porous medium equations

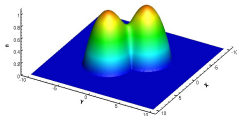
$$u_t - \Delta(u^2) = 0$$



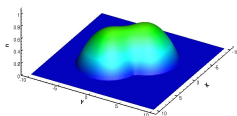
(a) $t = 0$



(b) $t = 0.1$



(c) $t = 0.5$



(d) $t = 2$

2-D porous medium equations

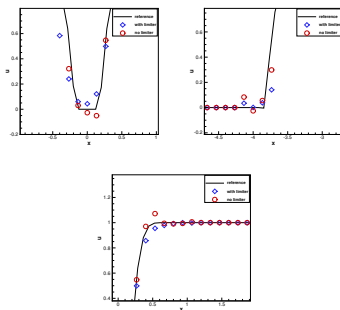
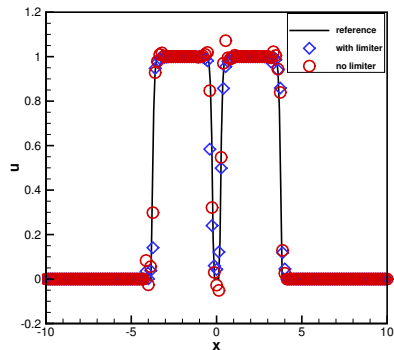


Figure: The cut of the DDG interface correction solution along line $x + y = 0$ at $t = 0.005$. Red circle symbol: no M-P-S limiter. Blue diamond symbol: M-P-S limiter applied.

2-D strongly degenerate nonlinear convection diffusion equations

$$u_t + \nabla \cdot \langle u^2, u^2 \rangle = \epsilon \nabla \cdot (\nu(u) \nabla u), \quad \nu(u) = \begin{cases} 0, & |u| \leq 0.25, \\ 1, & |u| > 0.25. \end{cases}$$